

# Finite Element Methods for a class of Stochastic PDEs from phase transitions

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## Phase separation

- In a vessel  $\mathcal{D}$  a two-phases mixture is forced to homogenization,
  - ↪ reducing the temperature the binary alloy starts to separate rapidly,
  - ↪ thin transition layers of width an order of  $\varepsilon$  are formed and the evolution slows down ( $\varepsilon$  small enough).
  - ↪  $\varepsilon$ -dependent spdes for the concentration.
  - ↪ sharp interface limit as  $\varepsilon \rightarrow 0$ .

### Some essentials for the numerical approximation:

- underlying existence-regularity theory,
- noise definition and its numerical implementation,
- sufficient theory (at least numerical convergence) so that the applied scheme will lead somewhere we provisionally trust.

## A class of $\varepsilon$ -dependent spdes

$$\begin{aligned} u_t &= -\delta(\varepsilon) \Delta(\varepsilon^2 \Delta u - f(u)) + \mu(\varepsilon)(\varepsilon^2 \Delta u - f(u)) + \mathcal{F}(u, x, t; \varepsilon), \\ x &\in \mathcal{D}, \quad t > 0, \\ u(x, 0) &= u_0(x; \varepsilon), \quad x \in \mathcal{D} \quad (\text{possibly layered}), \\ \text{and b.c. on } \partial\mathcal{D}, \end{aligned} \tag{1}$$

for  $u$  the concentration of one of the components of the **binary** alloy, **and**

$$0 < \varepsilon \ll 1.$$

$f(u) = F'(u)$ , is the derivative of a **double** equal-well potential, typical example

$$f(u) = u^3 - u, \quad F(u) := \frac{1}{4}(u^2 - 1)^2,$$

and the  $\varepsilon$ -dependent **weights** satisfy

$$\delta(\varepsilon) \geq 0, \quad \mu(\varepsilon) \geq 0.$$

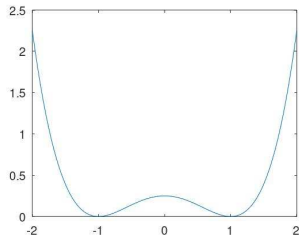


Figure: The potential  $F$

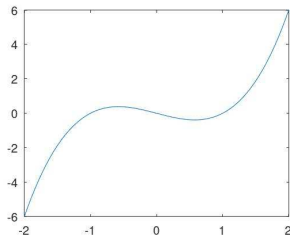


Figure: The nonlinearity  $f = F'$

## Main cases

- $\delta(\varepsilon) > 0$ ,  $\mu(\varepsilon) = 0$  Cahn-Hilliard equation (Model B).
- $\delta(\varepsilon) = 0$ ,  $\mu(\varepsilon) > 0$  Allen-Cahn equation (Model A).
- $\delta(\varepsilon) > 0$ ,  $\mu(\varepsilon) > 0$  'competition'  
of Cahn-Hilliard and Allen-Cahn operators (Model B/A).

$\mathcal{F} \equiv 0$  the homogeneous problem,

$$\mathcal{F} \neq 0 = \begin{cases} \text{smooth stochastic perturbation,} \\ \text{non smooth stochastic perturbation.} \end{cases}$$

**Questions:** Existence, regularity, transitions profile, sharp interface limit as  $\varepsilon \rightarrow 0$ , numerical approximation.

**Numerical Difficulty:**  $\varepsilon \ll 1$  in physical scale

$\hookrightarrow$  severe rounding errors.

## 1. The stochastic Cahn-Hilliard equation

$$\begin{aligned} u_t &= \Delta(-\varepsilon \Delta u + \varepsilon^{-1} f(u)) + \varepsilon^\gamma \dot{W}(x, t), \quad x \text{ in } \mathcal{D}, \\ \frac{\partial u}{\partial \nu} &= \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \partial \mathcal{D} \quad (\text{Neumann b.c.}) \end{aligned} \quad (2)$$

Here

$$\delta(\varepsilon) := \varepsilon^{-1}, \text{ and } \mu(\varepsilon) \equiv 0,$$

with

$$\mathcal{F}(u, x, t; \varepsilon) := \mathcal{F}(x, t; \varepsilon) = \varepsilon^\gamma \dot{W} = \varepsilon^\gamma \sum_{i=1}^{\infty} a_i \dot{\beta}_i(t) e_i(x),$$

for  $\beta_i$  stochastically independent brownian motions.

Let the chemical potential  $v = v(\varepsilon)$  and derive the stochastic system:

$$\begin{aligned} u_t &= -\Delta v + \varepsilon^\gamma \dot{W} \\ v &= -\varepsilon^{-1} f(u) + \varepsilon \Delta u. \end{aligned} \quad (3)$$

In  $d = 2, 3$ , by applying a boot-strap stochastic argument, in A, Blömker, Karali, A.I.H.P. Prob.-Stat, 2018, we proved the next theorem for the Stochastic sharp interface limit.

**Theorem:** Let  $\gamma > \gamma_0 > 1$ . The **limit** of  $v$  as  $\varepsilon \rightarrow 0^+$  solves on  $[0, T]$  the **deterministic Hele-Shaw problem**

$$\Delta v = 0 \quad \text{in } \mathcal{D} \setminus \Gamma(t), \quad t > 0$$

$$\partial_n v = 0 \quad \text{on } \partial \mathcal{D}$$

$$v = \lambda H \quad \text{on } \Gamma(t)$$

$$V = \frac{1}{2}(\partial_n v^+ - \partial_n v^-) \quad \text{on } \Gamma(t)$$

$$\Gamma(0) = \Gamma_0,$$

where  $H$  is the mean curvature, and  $V$  the velocity of the surface  $\Gamma(t)$ .

**Remarks:**

- The above was derived when

$$\int_{\mathcal{D}} W(x, t) dx = 0 \quad (\text{mass conservation}).$$

- This problem **coincides** to the limit problem of deterministic C-H of Alikakos, Bates, Chen, ARMA 1994.

- When  $\gamma$  **smaller** formal asymptotics indicated convergence but with

$$v = \lambda H + \mathbf{W} \quad \text{on } \Gamma(t).$$

## Numerical approximation

In (2), we consider **one mode** in the noise series, i.e.,

$$\mathcal{F}(u, x, t; \varepsilon) := \mathcal{F}(x, t; \varepsilon) = \varepsilon^\gamma \dot{W} = \varepsilon^\gamma \sigma(x) \dot{\beta}(t).$$

The **non-smooth** in time noise is approximated, by using the increments of the brownian motion  $\beta(t) \equiv N(0, t)$ , as follows

$$\dot{\beta}(t) \cong \text{rate of change} = \frac{\beta(t^i) - \beta(t^{i-1})}{t^i - t^{i-1}} \sim k^{-1/2} N(0, 1) = N(0, k^{-1}).$$

In A, Banas, Nürnberg, Prohl, Numer. Math. 2021, we defined the next **semi-discretization** in time:

**Let**  $0 = t^0 < t^1 < \dots < t^J = T$  a uniform partition, with step-size  $k = \frac{T}{J}$ . For  $1 \leq j \leq J$ , let

$$\Delta_j \beta := \beta(t^j) - \beta(t^{j-1}) \sim N(0, k),$$

the corresponding **brownian increments**.



We seek in  $H^1(\mathcal{D})$  approximations of the solution  $u$  and of the chemical potential  $v$

$$U^j(x) \simeq u(x, t^j), \quad V^j(x) \simeq v(x, t^j)$$

satisfying the **discrete** weak formulation of the **system** (3)

$$(U^j - U^{j-1}, \varphi) + k(\nabla W^j, \nabla \varphi) = \varepsilon^\gamma(\sigma, \varphi) \Delta_j \beta \quad \forall \varphi \in H^1(\mathcal{D}),$$

$$\varepsilon(\nabla U^j, \nabla \psi) + \frac{1}{\varepsilon}(f(U^j), \psi) = (V^j, \psi) \quad \forall \psi \in H^1(\mathcal{D}),$$

$$U^0(x) = u_0(x; \varepsilon) \in H^1(\mathcal{D}).$$

**Remark:** The scheme without **noise** was introduced and analyzed in Feng, Prohl Numer. Math. 2004.

In dimensions  $d = 2$ ,

- by a **discrete** boot-strap argument, we estimated the error in expectation in  $H^{-1}$  for  $k < \varepsilon^\ell$  for some  $\ell > 0$  sufficiently **large**,
- the **convergence** to the Hele-Shaw was **experimentally** investigated for a fully discrete space-time scheme.

**Remark:** This result was derived by energy estimates, i.e., up to  $H^1(\mathcal{D})$  regularity used for the stochastic C-H.

But how smooth the solution can really be?

Let us consider the more general case of **multiplicative** noise, where

$$\mathcal{F}(u, x, t; \varepsilon) := \sigma(u) \varepsilon^\gamma \dot{W} = \sigma(u) \varepsilon^\gamma \sum_{i=1}^{\infty} a_i \dot{\beta}_i(t) e_i(x).$$

In A Nonlinearity, 2023, the next result was derived.

**Theorem:** Let  $\|u_0\|_{H^2(\mathcal{D})}$  have bounded  $p$ -moments,  $d = 1, 2, 3$ , and

$$|\sigma(x)| + |\sigma'(x)| + |\sigma''(x)| \leq c, \quad \forall x \in \mathbb{R}, \quad \text{and} \quad \sum_{i=1}^{\infty} a_i^2 \|e_i\|_{H^4(\mathcal{D})}^2 < \infty.$$

**Then** for any  $T > 0$ , and some  $k = k(p) > 0$ , and for any  $p \geq 1$

$$\mathbb{E} \left( \sup_{0 \leq t \leq T} \|u\|_{L^\infty(\mathcal{D})}^{2p} \right) \leq c \mathbb{E} \left( \sup_{0 \leq t \leq T} \|u\|_{H^2(\mathcal{D})}^{2p} \right) \leq c \varepsilon^{-k(p)}.$$

Moreover, if  $\partial\mathcal{D}$  is  $C^1$ , for

$$\|u(\cdot, t)\|_{C^{0,\theta}(\overline{\mathcal{D}})} := \sup_{x \neq y \in \overline{\mathcal{D}}} \frac{|u(x, t) - u(y, t)|}{|x - y|^\theta},$$

it follows for any  $0 < \theta < 1$  for  $d = 2$ , and for  $\theta = \frac{1}{2}$  when  $d = 3$

$$\mathbb{E} \left( \sup_{0 \leq t \leq T} \|u(\cdot, t)\|_{C^{0,\theta}(\overline{\mathcal{D}})}^{2p} \right) \leq c \mathbb{E} \left( \sup_{0 \leq t \leq T} \|u(\cdot, t)\|_{H^2(\mathcal{D})}^{2p} \right) \leq c \varepsilon^{-k(p)}.$$

Additionally, if the initial data are **layered**, then  $u \rightarrow \pm 1$  a.s. In particular, for any  $\gamma > 0$  (of noise), and any  $0 < \alpha < \min\{\gamma, 1/2\}$

$$\lim_{\varepsilon \rightarrow 0^+} P(\| |u| - 1 \|_{L^2(\mathcal{D})} \geq \varepsilon^\alpha) = 0.$$

**Remark:** The  $H^2(\mathcal{D})$  estimate, establishes a.s. convergence in  $L^2(\mathcal{D})$  to the **step function** through a.s. continuous paths.

The  $H^2$  regularity was derived by applying a proper version of **BDG inequality** (see for example in Marinelli, Röckner, Exp. Math. 2016), for the **a priori estimates** in expectation up to  $H^2$  (inspired by Elliott, Songmu, ARMA, 1986, made for the determ. CH.)

## 2. The stochastic Allen-Cahn equation with mild noise

$$\begin{aligned}w_t &= \Delta w + \frac{f(w)}{\varepsilon^2} + \frac{\dot{W}(x, t; \varepsilon)}{\varepsilon}, \quad x \in \Omega, \quad 0 < t \leq T, \\w(x, 0) &= w_0(x), \quad x \in \Omega, \\ \frac{\partial w}{\partial \eta} &= 0, \quad x \in \partial\Omega, \quad 0 < t \leq T \quad (\text{Neumann b.c.})\end{aligned}\tag{4}$$

Here

$$\delta(\varepsilon) \equiv 0, \quad \text{and} \quad \mu(\varepsilon) = \varepsilon^{-2}.$$

with

$$\mathcal{F}(u, x, t; \varepsilon) := \mathcal{F}(x, t; \varepsilon) = \frac{\dot{W}(x, t; \varepsilon)}{\varepsilon},$$

for

$\dot{W}(x, t; \varepsilon)$  a mild noise tending to rough as  $\varepsilon \rightarrow 0$ .)

- In A IMA NA 2020, a **DG nonlinear** method was introduced

↪ **existence** and **uniqueness** of the numerical solution were proved

↪ optimal **error** estimates were derived.

- In A, Egwu, Yan IMA NA 2023, the **mild noise** was numerically approximated

↪ and we applied an **a posteriori** error analysis on the scheme.

**Smooth time-noise definition** (H. Weber AIHP P.Stat., 2010)

Let  $\rho : \mathbb{R} \rightarrow \mathbb{R}^+$

- compactly supported in  $[-1, 1]$ ,
- symmetric around zero,
- it satisfies

$$\int_{-\infty}^{\infty} \rho(s) ds = 1.$$

Define for  $0 < \gamma < \frac{2}{3}$

$$\rho^\varepsilon(s) := \varepsilon^{-\gamma} \rho\left(\frac{s}{\varepsilon^\gamma}\right).$$

Let  $\beta(t), \tilde{\beta}(t) \sim N(0, t)$ ,  $t \geq 0$  **two stochastically independent Brownian Motions**.

The **smooth approximation of a brownian**  $\beta(t)$  is given by,

$$\beta_\varepsilon(t) := \int_{-\infty}^0 \rho^\varepsilon(t-s) \tilde{\beta}(-s) ds + \int_0^\infty \rho^\varepsilon(t-s) \beta(s) ds,$$

and the smooth approximation of its formal derivative by

$$\dot{\beta}_\varepsilon(t) := \int_{-\infty}^0 \partial_t(\rho^\varepsilon(t-s)) \tilde{\beta}(-s) ds + \int_0^\infty \partial_t(\rho^\varepsilon(t-s)) \beta(s) ds,$$

which is the smooth approximation of  $\dot{\beta}$ .

## Numerical implementation of smooth noise

How we insert  $x$ -dependence in the mild noise?

An answer: In A IMA NA 2020, the smooth space-time noise was defined by the convolution

$$\dot{W}(x, t; \varepsilon) := G(x) \dot{\beta}_\varepsilon(t),$$

for  $G \in C^\infty$ .

How we select  $\rho$ ?

How we construct the convolution integral numerically?

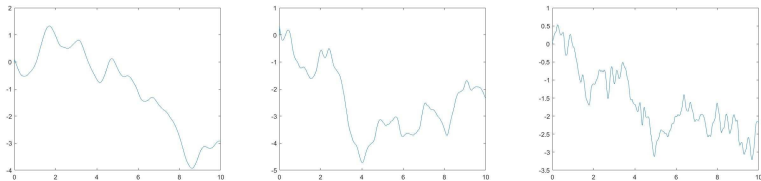
In A, Egwu, Yan IMA NA 2023, the bump function was selected and the composite trapezoidal rule on the convolution was used.

Let the compactly supported **bump function**  $r : \mathbb{R} \rightarrow \mathbb{R}^+$ , with

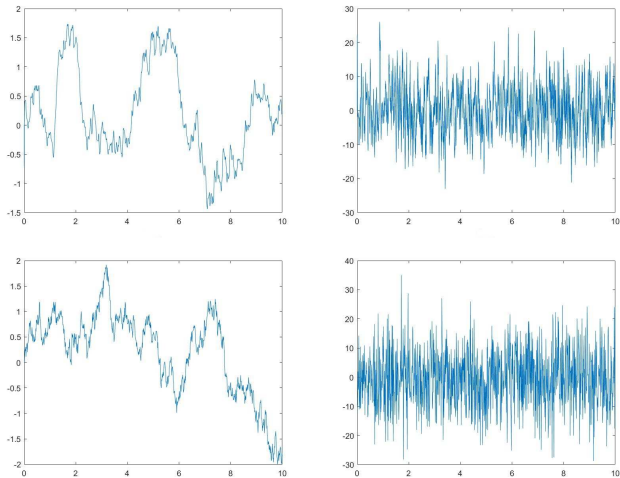
$$r(s) := \begin{cases} e^{-\frac{1}{1-s^2}} & s \in (-1, 1) \\ 0 & \text{otherwise} \end{cases}, \text{ and define}$$

$$\rho(s) := r(s) \left[ \int_{-1}^1 r(s) ds \right]^{-1} = \begin{cases} e^{-\frac{1}{1-s^2}} \left[ \int_{-1}^1 e^{-\frac{1}{1-s^2}} ds \right]^{-1} & s \in (-1, 1), \\ 0 & \text{otherwise.} \end{cases}$$

Then  $\rho$  satisfies all the properties required in the convolution definitions.



**Figure:** The smooth **approximation**  $\beta_\epsilon(t)$  of the brownian for  $\gamma = \frac{1}{8} < \frac{2}{3}$  and  $T = 10$ , for  $\epsilon = 10^{-3}, 10^{-6}, 10^{-9}$  (from the left to the right).



**Figure:** First row: The **smooth** approximation  $\beta_\varepsilon(t)$  of the **brownian** and  $\dot{\beta}_\varepsilon(t)$  the **smooth** approximation of the **formal derivative** of the brownian, for  $\gamma = \frac{1}{8} < \frac{2}{3}$ ,  $T = 10$ , and  $\varepsilon = 10^{-14}$ . Second row: A **brownian** path (with D. Higham's algorithm, SIAM Rev. 2001), and its **formal derivative**.



The DG scheme for the mild stochastic AC,  
motivated by Jamet SINUM 1978.

We **apply** the exponential transformation

$$w = e^{b(\varepsilon)t} u, \quad \text{with} \quad \inf_{\varepsilon \in (0,1)} (b(\varepsilon) - \varepsilon^{-2}) \geq \hat{c}_0 > 0,$$

which is analogous to these of A, Plexousakis Numer.Math. 2010,  
A, Plexousakis ESAIM M2AN 2019, and **obtain**

$$\begin{aligned} u_t = & \Delta u - b(\varepsilon)u + \frac{g(u, \varepsilon, t)}{\varepsilon^2} \\ & + m(\varepsilon, t) \frac{\dot{W}(x, t; \varepsilon)}{\varepsilon}, \quad x \in \mathcal{D}, \quad 0 < t \leq T, \end{aligned}$$

**for**

$$g(u, \varepsilon, t) := u - e^{2b(\varepsilon)t} u^3, \quad m(\varepsilon, t) := e^{-b(\varepsilon)t}.$$

The exponential transformation resulted to **numerical stability** on:

- linear parabolic equations (linear Schrödinger, Heat)
- stochastic Allen-Cahn.

**Let**  $0 = t^0 < t^1 < \dots < t^N = T$ , a partition of  $[0, T]$ , and set

$$G^n := \mathcal{D} \times (t^n, t^{n+1}).$$

For each  $0 \leq n \leq N - 1$

let a family  $\{V_h^n\}$  of finite dimensional subspaces of  $H^1(G^n)$ , and **define**  $V_h$  such that

$$V_h|_{(t^n, t^{n+1})} := V_h^n.$$

**Remark:** The functions of  $V_h$  are in general **discontinuous** at the temporal nodes  $t^n$ .

↪ We seek  $u_h \in V_h$  such that:

$$\begin{aligned} & - ((u_h, \partial_t v_h))_{G^n} + ((\nabla u_h, \nabla v_h))_{G^n} + b(\varepsilon)((u_h, v_h))_{G^n} \\ & - \varepsilon^{-2}((u_h, v_h))_{G^n} + \varepsilon^{-2}((e^{2b(\varepsilon)t}(u_h)^3, v_h))_{G^n} \\ & + (u_h^{n+1}, v_h^{n+1})_{\mathcal{D}} - (u_h^n, v_h^{n+0})_{\mathcal{D}} = \varepsilon^{-1}((e^{-b(\varepsilon)t}\xi_t^\varepsilon, v_h))_{G^n}, \\ & \forall v_h \in V_h^n, \quad n = 0, \dots, N-1, \quad u_h^0 = u_0, \end{aligned}$$

for  $\xi_t^\varepsilon := \dot{W}(x, t; \varepsilon)$  the mild noise.

Here,  $((\cdot, \cdot))_{G^n}$ ,  $(\cdot, \cdot)_{\mathcal{D}}$  denote the  $L^2(G^n)$  and the  $L^2(\mathcal{D})$  inner products respectively, while  $|\cdot|_{\mathcal{D}} := \|\cdot\|_{L^2(\mathcal{D})}$ .

**Remarks:** The DG scheme is **nonlinear**, **time** discontinuous, **adaptive** in  $t$ , while  $u_h^0$  coincides with the continuous problem initial data.

The **weak** formulation and the inner product is defined by  $L^2$  **double** integration in space-time.

### 3. Coexistence of both Cahn-Hilliard and Allen-Cahn operators

Let us recall the CH/AC equation

$$u_t = -\delta(\varepsilon)\Delta(\varepsilon^2\Delta u - f(u)) + \mu(\varepsilon)(\varepsilon^2\Delta u - f(u)) + \mathcal{F}, \quad x \in \mathcal{D}, \quad t > 0,$$

where

$$\delta(\varepsilon) > 0, \quad \mu(\varepsilon) \geq 0.$$

For the homogeneous problem ( $\mathcal{F} \equiv 0$  and so **deterministic**)  
in  $d = 1$ , A, Karali, Tzirakis, in Cal.Var. 2021

↪ we estimated the **spectrum** of the linearized CH/AC operator,

↪ determined **families of weights**  $\delta(\varepsilon)$ ,  $\mu(\varepsilon)$ , for which the  
dynamics of  $N$  layers are **stable** and **rest** exponentially small in  $\varepsilon$ ,

↪ derived a system of ODEs for the layers **motion**.

For the **same** problem, in A, Karali, D. Li, J.Sc.Comput. 2025,

↪ we applied a **SAV**-linearized RK method in  $d = 1, 2, 3$

- we proved optimal **error** estimates,
- we had the first **simulations** for the mixed problem (deterministic) by further discretizing in space by **FEM**.

**In progress:** The **SAV** method for **stochastic** equations.

Thank you!

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