

# High-order sampling of the invariant distribution of ergodic stochastic dynamics: preconditioning and postprocessing

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based on joint works with

Eugen Bronasco (Göteborg), Benedict Leimkuhler, and Dominic Phillips (Edinburgh)

Milstein's method: 50 years on, Nottingham, June 2025

## Long time accuracy for ergodic stochastic problems

$$dX(t) = f(X(t))dt + g(X(t))dW(t), \quad X(0) = x.$$

Under standard **ergodicity assumptions**,

$$\begin{aligned} \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \phi(X(t)) dt &= \int_{\mathbb{R}^d} \phi(y) d\mu_\infty(y) \\ \left| \mathbb{E}(\phi(X(t))) - \int_{\mathbb{R}^d} \phi(y) d\mu_\infty(y) \right| &\leq K(x, \phi) e^{-ct}, \quad \text{for all } t \geq 0. \end{aligned}$$

Different types of convergence as  $h \rightarrow 0$ :

- **strong order  $q$**  (fixed finite time  $T$ ),

$$\mathbb{E}(|X(t_n) - X_n|) \leq Ch^q, \quad \text{for all } t_n = nh \leq T.$$

- **weak order  $r$**  (fixed finite time  $T$ ),

$$|\mathbb{E}(\phi(X(t_n))) - \mathbb{E}(\phi(X_n))| \leq Ch^r, \quad \text{for all } t_n = nh \leq T.$$

- **order  $p$  for the invariant measure** (long time behavior): in general  $p \geq r \geq q$ .

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Two standard approaches using an ergodic integrator of **order  $p$** :

- Compute a single long trajectory  $\{X_n\}$  of length  $T = Nh$ ,

$$\frac{1}{N+1} \sum_{k=0}^N \phi(X_k) \simeq \int_{\mathbb{R}^d} \phi(y) d\mu_\infty(y), \quad \text{error } \mathcal{O}(h^p + T^{-1/2}),$$

- Compute many trajectories  $\{X_n^i\}$  of length  $t = Nh$ ,

$$\frac{1}{M} \sum_{i=1}^M \phi(X_N^i) \simeq \int_{\mathbb{R}^d} \phi(y) d\mu_\infty(y), \quad \text{error } \mathcal{O}(e^{-ct} + h^p + M^{-1/2}).$$

## Example: Overdamped Langevin equation (Brownian dynamics)

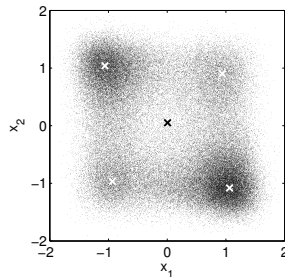
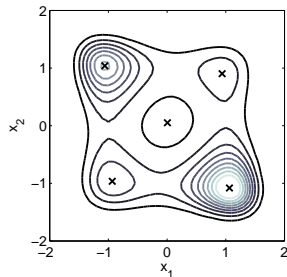
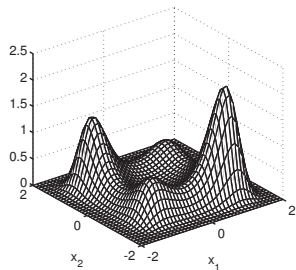
$$dX(t) = -\nabla V(X(t))dt + \sqrt{2}dW(t).$$

$X(t)$ : solution process in  $\mathbb{R}^d$ .  $W(t)$ : standard Brownian motion in  $\mathbb{R}^d$ .

Ergodicity: invariant measure  $\mu_\infty$  has Gibbs density  $\rho_\infty(x) = Ze^{-V(x)}$ ,

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \phi(X(s))ds = \int_{\mathbb{R}^d} \phi(y)d\mu_\infty(x), \quad a.s.$$

Example ( $d = 2$ ):  $V(x) = (1 - x_1^2)^2 + (1 - x_2^2)^2 + \frac{x_1 x_2}{2} + \frac{x_2}{5}$ .





## Example: Overdamped Langevin equation (Brownian dynamics)

$$dX(t) = -\nabla V(X(t))dt + \sqrt{2}dW(t).$$

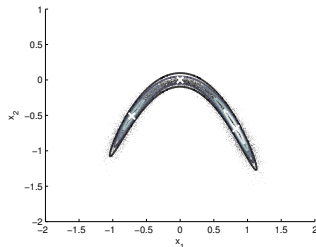
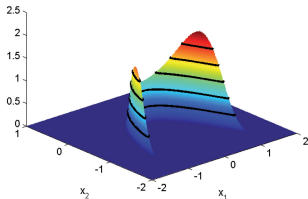
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Example ( $d = 2$ ), Stiff case:

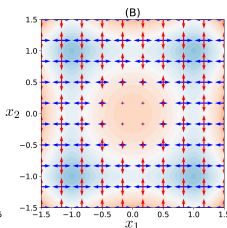
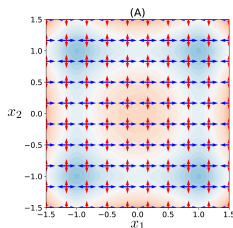
$$V(x) = (1 - x_1^2)^2 + x_2^4 - x + x_3 \cos(x_2) + 100(x_2 + x_1^2)^2 + \frac{10^6}{2}(x_1 - x_3)^2.$$



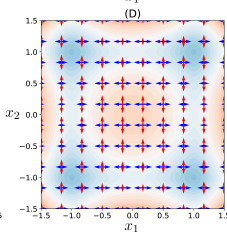
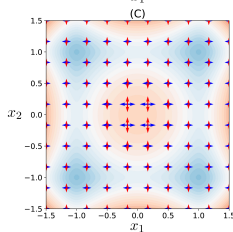
# Variable diffusion case: preconditioning the ergodic dynamics

$$dX = (-D^T D(X) \nabla V(x) + \text{div}(D^T D)(X)) dt + \sqrt{2} D(X) dW, \quad (\text{variable } D(x) \in \mathbb{R}^{d \times d})$$

$$D_A(x) = \begin{pmatrix} 2 & 0 \\ 0 & 1.5 \end{pmatrix}$$



$$D_C(x) = (1 + 5e^{-\|x\|^2/0.18})I$$



$$D_B(x) = \frac{1}{1+5e^{-\|x\|^2/0.18}} I$$

$$D_D(x) = I - \frac{xx^T}{2\|x\|^2+1}$$

$$\text{Quadruple well potential } V(x) = \sqrt{\frac{17}{16} - 2x_1^2 + x_1^4} + \sqrt{\frac{17}{16} - 2x_2^2 + x_2^4}.$$

## Variable diffusion case: preconditioning the ergodic dynamics

Consider the Langevin dynamics with smooth variable diffusion  $D(x) \in \mathbb{R}^{d \times d}$  symmetric and uniformly positive definite (same invariant measure with  $\rho_\infty(x) = Ze^{-\frac{\sigma^2}{2}V(x)}$ ),

$$dX = F(X)dt + \sigma D(X)dW, \quad F(x) = -D^2(x)\nabla V(x) + \frac{\sigma^2}{2}\operatorname{div}(D^2)(x).$$

The Leimkuhler-Matthews (2013) scheme:

$$X_{n+1} = X_n + hF(X_n) + \sigma D \frac{\xi_n + \xi_{n+1}}{2}, \quad \xi_n \sim \mathcal{N}(0, I) \text{ i.i.d.}$$

**Challenge:** generalization of the Leimkuhler-Matthews in the variable diffusion case?

- For  $D(x) = I$  (identity matrix), it yields order two for the invariant measure.
- For  $D(x) = D$  (constant matrix), it yields order two for the invariant measure.
- For  $D(x)$ , variable, it is **not consistent** in general. The modified version

$$X_{n+1} = X_n + hF(X_n) + \frac{h\sigma^2}{4} \operatorname{div}(D^2)(x) + \sigma D \frac{\xi_n + \xi_{n+1}}{2},$$

has **order one** in the one-dimensional case  $d = 1$ , but is **not consistent** in general.

# Plan of the talk

- 1 Part 1: High order for the invariant measure and postprocessors
  - Weak order condition and order conditions sampling for the invariant measure
  - Postprocessors for sampling for the invariant measure
- 2 Part 2: High-order in the variable diffusion case
  - New generalization of the Leimkuhler-Matthews Scheme
  - Bias analysis based on exotic aromatic Butcher-series

## 1 Part 1: High order for the invariant measure and postprocessors

- Weak order condition and order conditions sampling for the invariant measure
- Postprocessors for sampling for the invariant measure

- Abdulle, V., Zygalakis, *High order numerical approximation of ergodic SDE invariant measures*, [SIAM SINUM](#), 2014.
- Abdulle, V., Zygalakis, *Long time accuracy of Lie-Trotter splitting methods for Langevin dynamics*, [SIAM SINUM](#), 2015.
- G. V., *Postprocessed integrators for the high order integration of ergodic SDEs*, [SIAM SISC](#), 2015.

## A classical tool: the Fokker-Plank equation

$$dX(t) = f(X(t))dt + \sqrt{2}dW(t).$$

The density  $\rho(x, t)$  of  $X(t)$  at time  $t$  solves the parabolic problem

$$\partial_t \rho = \mathcal{L}^* \rho = -\operatorname{div}(f \rho) + \Delta \rho, \quad t > 0, x \in \mathbb{R}^d.$$

For ergodic SDEs, for any initial condition  $X(0) = X_0$ , as  $t \rightarrow +\infty$ ,

$$\mathbb{E}(\phi(X(t))) = \int_{\mathbb{R}^d} \phi(x) \rho(x, t) dx \longrightarrow \int_{\mathbb{R}^d} \phi(x) d\mu_\infty(x).$$

The invariant measure  $d\mu_\infty(x) \sim \rho_\infty(x)dx$  is a stationary solution ( $\partial_t \rho_\infty = 0$ ) of the Fokker-Plank equation

$$\mathcal{L}^* \rho_\infty = 0.$$

# Asymptotic expansions

Theorem (Talay and Tubaro, 1990, see also, Milstein, Tretyakov)

Assume that  $X_n \mapsto X_{n+1}$  (weak order  $p$ ) is *ergodic* and has a Taylor expansion

$$\mathbb{E}(\phi(X_1)|X_0 = x) = \phi(x) + h\mathcal{L}\phi + h^2 A_1\phi + h^3 A_2\phi + \dots$$

If  $\mu_\infty^h$  denotes the numerical invariant distribution, then

$$e(\phi, h) = \int_{\mathbb{R}^d} \phi d\mu_\infty^h - \int_{\mathbb{R}^d} \phi d\mu_\infty = \lambda_p h^p + \mathcal{O}(h^{p+1}),$$

$$\mathbb{E}(\phi(X_n)) - \int_{\mathbb{R}^d} \phi d\mu_\infty - \lambda_p h^p = \mathcal{O}(\exp(-cnh) + h^{p+1}),$$

where, denoting  $u(t, x) = \mathbb{E}\phi(X(t, x))$ ,

$$\begin{aligned} \lambda_p &= \int_0^{+\infty} \int_{\mathbb{R}^d} \left( A_p - \frac{\mathcal{L}^{p+1}}{(p+1)!} \right) u(t, x) \rho_\infty(x) dx dt \\ &= \int_0^{+\infty} \int_{\mathbb{R}^d} u(t, x) (A_p)^* \rho_\infty(x) dx dt. \end{aligned}$$

## High order approximation of the numerical invariant measure

Assume that  $X_n \mapsto X_{n+1}$  is **ergodic** with standard assumptions and

$$\mathbb{E}(\phi(X_1)|X_0 = x) = \phi(x) + h\mathcal{L}\phi + h^2 A_1 \phi + h^3 A_2 \phi + \dots$$

### Standard weak order condition.

If  $A_j = \frac{\mathcal{L}^j}{j!}$ ,  $1 \leq j < p$ , then (weak order  $p$ )

$$\mathbb{E}(\phi(X(t_n))) = \mathbb{E}(\phi(X_n)) + \mathcal{O}(h^p), \quad t_n = nh \leq T.$$

### Order condition for the invariant measure.

If  $A_j^* \rho_\infty = 0$ ,  $1 \leq j < p$ , then (order  $p$  for the invariant measure)

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N \phi(X_n) = \int_{\mathbb{R}^d} \phi(y) d\mu(y) + \mathcal{O}(h^p),$$
$$\mathbb{E}(\phi(X_n)) - \int_{\mathbb{R}^d} \phi d\mu_\infty = \mathcal{O}(\exp(-cnh) + h^p).$$







## Order conditions for the invariant measure sampling

$$Y_i^n = X_n + h \sum_{j=1}^s a_{ij} f(Y_j^n) + d_i \sigma \sqrt{h} \xi_n, \quad i = 1, \dots, s,$$

$$X_{n+1} = X_n + h \sum_{i=1}^s b_i f(Y_i^n) + \sigma \sqrt{h} \xi_n,$$

Theorem (Laurent, V., 2020, Conditions for order  $p$  for the invariant measure)

Order	Tree $\tau$	$F(\tau)(\phi)$	Order condition
1		$\phi' f$	$\sum b_i = 1$
2		$\phi' f' f$	$\sum b_i c_i - 2 \sum b_i d_i = -\frac{1}{2}$
		$\phi' \Delta f$	$\sum b_i d_i^2 - 2 \sum b_i d_i = -\frac{1}{2}$
3		$\phi' f' f' f$	$\sum b_i a_{ij} c_j - 2 \sum b_i a_{ij} d_j$ $+ \sum b_i c_i - (\sum b_i d_i)^2 = 0$
	...	...	...

# Postprocessed integrators for ergodic SDEs

Idea: extend to the context of ergodic SDEs the popular idea of **effective order** for ODEs from Butcher 69',

$$y_{n+1} = \chi_h \circ K_h \circ \chi_h^{-1}(y_n), \quad y_n = \chi_h \circ K_h^n \circ \chi_h^{-1}(y_0).$$

## Example based on the Euler-Maruyama method

for Brownian dynamics:  $dX(t) = -\nabla V(X(t))dt + \sigma dW(t)$ .

$$X_{n+1} = X_n - h\nabla V\left(X_n + \frac{1}{2}\sigma\sqrt{h}\xi_n\right) + \sigma\sqrt{h}\xi_n, \quad \bar{X}_n = X_n + \frac{1}{2}\sigma\sqrt{h}\xi_n.$$

$X_n$  has order 1 of accuracy for the invariant measure.

$\bar{X}_n$  has order 2 of accuracy for the invariant measure (postprocessor).

First derived as a **non-Markovian method** by Leimkuhler, Matthews (2013), see Leimkuhler, Matthews, Tretyakov (2014),

$$\bar{X}_{n+1} = \bar{X}_n - h\nabla V(\bar{X}_n) + \frac{1}{2}\sigma\sqrt{h}(\xi_n + \xi_{n+1}).$$

## Postprocessed integrators for ergodic SDEs: nonlinear case

Postprocessing:  $\bar{X}_n = G_n(X_n)$ , with weak Taylor series expansion

$$\mathbb{E}(\phi(G_n(x))) = \phi(x) + h^p \bar{A}_p \phi(x) + \mathcal{O}(h^{p+1}).$$

### Theorem (V., 2015)

*Under technical assumptions, assume that  $X_n \mapsto X_{n+1}$  and  $\bar{X}_n$  satisfy*

$$A_j^* \rho_\infty = 0 \quad j < p, \quad (\text{order } p \text{ for the invariant measure}),$$

*and*

$$(A_p + [\mathcal{L}, \bar{A}_p])^* \rho_\infty = (A_p + \mathcal{L} \bar{A}_p - \bar{A}_p \mathcal{L})^* \rho_\infty = 0,$$

*then (order  $p + 1$  for the invariant measure)*

$$\mathbb{E}(\phi(\bar{X}_n)) - \int_{\mathbb{R}^d} \phi d\mu_\infty = \mathcal{O}(\exp(-cnh) + h^{p+1}).$$

**Remark:** postprocessing is needed only at the end of the time interval (not each step).

**Related:** Bréhier, V., SIAM SISC 2016, **order improved by +1** for the stochastic heat eq.

## Part 2: High-order in the variable diffusion case

### 1 Part 1: High order for the invariant measure and postprocessors

- Weak order condition and order conditions sampling for the invariant measure
- Postprocessors for sampling for the invariant measure

### 2 Part 2: High-order in the variable diffusion case

- New generalization of the Leimkuhler-Matthews Scheme
- Bias analysis based on exotic aromatic Butcher-series

- Laurent, V., *Exotic aromatic B-series for the study of long time integrators for a class of ergodic SDEs*, *Math. Comp.*, 2020.
- Laurent, V., *Order conditions for sampling the invariant measure of ergodic stochastic differential equations on manifolds*, *FoCM*, 2022.
- Bronaco, Leimkuhler, Phillips, V.. Order two scheme for the invariant measure sampling of Langevin dynamics with variable diffusion, submitted, 2025.

# New generalization of the Leimkuhler-Matthews Scheme

Consider the Langevin dynamics with smooth variable diffusion  $D(x) \in \mathbb{R}^{d \times d}$  symmetric and uniformly positive definite (same invariant measure with  $\rho_\infty(x) = Ze^{-\frac{2}{\sigma}V(x)}$ ),

$$dX = F(X)dt + \sigma D(X)dW, \quad F(x) = -D^2(x)\nabla V(x) + \frac{\sigma^2}{2}\operatorname{div}(D^2)(x).$$

The **new postprocessed scheme** has the form

$$\begin{aligned} X_{n+1} &= X_n + hF(\bar{X}_n) + \hat{\Phi}_h^D\left(X_n + \frac{h}{4}F(\bar{X}_{n-1})\right), \\ \bar{X}_n &= X_n + \frac{1}{2}\sigma\sqrt{h}D(X_n)\xi_n, \end{aligned}$$

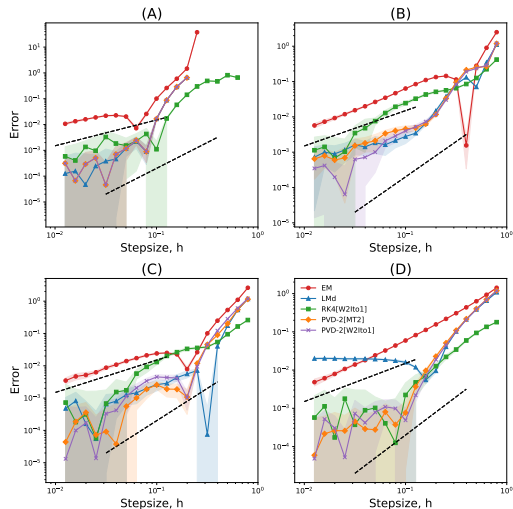
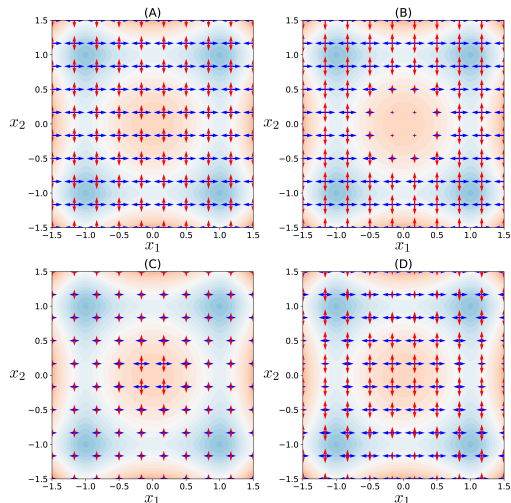
where  $I + \hat{\Phi}_h^D$  is a weak order 2 integrator of  $dX = \sigma D(X)dW$ .

**Theorem (Bronasco, Leimkuhler, Phillips, and V., 2025, submitted)**

The postprocessed method with  $\bar{X}_n$  above is second-order for sampling the invariant measure and has only one evaluation of  $\nabla V$  per step.

# Numerical experiments: Quadruple well potential

$$dX = (-D^2(X)\nabla V(x) + \text{div}(D^2)(X))dt + \sqrt{2}D(X)dW, \quad (\text{symmetric } D(x) \in \mathbb{R}^{d \times d})$$



# Aromatic Butcher-series

**Deterministic B-series:** Hairer & Wanner, 1972, using Butcher's seminal work (1960s).

Link with Hopf algebras of trees in quantum physics (Connes, Kreimer, 1980s).

**Stochastic case:** Tree formalism for **strong and weak errors on finite time**: Burrage K., Burrage P.M., 1996; Komori, Mitsui, Sugiura, 1997; Rößler, 2004/2006, ...

Here we focus of the accuracy for the **invariant measure (long time)**. We rewrite high-order differentials with trees. We denote  $F(\gamma)(f)$  the elementary differential of a tree  $\gamma$ .

For deterministic B-series:  $F(\bullet)(f) = f$ ,  $F(\bullet \mid \bullet)(f) = f'f$ ,  $F(\begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \end{array})(f) = f''(f, f'f)$

**Aromatic forests:** introduced for deterministic geometric integration for the study of **volume preservation** independently by Chartier, Murua, 2007, and Iserles, Quispel, Tse, 'K-loops' 2007 (see algebra structures in Bogfjellmo, 2015, Laurent, McLachlan, Munthe-Kaas, Verdier, 2024)

$$F(\begin{array}{c} \bullet \\ \circlearrowleft \\ \bullet \end{array} \mid \begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \end{array} \mid \begin{array}{c} \bullet \\ \bullet \end{array})(\phi) = \text{div}(f) \times \left( \sum \partial_i f_j \partial_j f_i \right) \times \phi' f$$

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Here we focus of the accuracy for the **invariant measure (long time)**. We rewrite high-order differentials with trees. We denote  $F(\gamma)(\phi)$  the elementary differential of a tree  $\gamma$ .

For S-series (Murua, 1999) :  $F(\bullet)(\phi) = \phi$ ,  $F(\bullet)(\phi) = \phi'f$ ,  $F(\begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \end{array})(\phi) = \phi''(f, f'f)$

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# Computing the expectation using lianas: examples



Grafted aromatic forests (like P-series):  $\xi \sim \mathcal{N}(0, I_d)$  is represented by crosses.

$$\mathbb{E}[F(\text{diagram})(\phi)] = \mathbb{E}[\phi'(f''(\xi, \xi))] = \sum_{i,j,k=1}^d \partial_i \phi \partial_{jk} f_i \mathbb{E}[\xi_j \xi_k] = \sum_{i,j=1}^d \partial_i \phi \partial_{jj} f_i = \phi' \Delta f = F(\text{diagram})(\phi),$$


$$\begin{aligned} \mathbb{E}[F(\text{diagram})(\phi)] &= \mathbb{E}[\phi^{(4)}(\xi, \xi, \xi, \xi)] = \sum_{i,j,k,l} \partial_{i,j,k,l} \phi \mathbb{E}[\xi_i \xi_j \xi_k \xi_l] \\ &= \sum_i \partial_{i,i,i,i} \phi \mathbb{E}[\xi_i^4] + 3 \sum_{\substack{i,j \\ i \neq j}} \partial_{i,i,j,j} \phi \mathbb{E}[\xi_i^2] \mathbb{E}[\xi_j^2] \\ &= 3 \sum_{i,j} \partial_{i,i,j,j} \phi = 3 \Delta^2 \phi = 3 F(\text{diagram})(\phi). \end{aligned}$$


Idea of the proof for general trees: L. Isserlis' theorem (Wick's probability theorem).


## New exotic aromatic B-series: using **lianas**

$$F(\text{diagram 1})(\phi) = \phi''(f'\xi, \xi) \quad \text{and} \quad F(\text{diagram 2})(\phi) = \phi' f''(\xi, \xi).$$



We introduce **lianas** in our trees, now called **exotic aromatic trees and forests**:

$$F(\text{diagram 3}) = \sum_i \phi''(f' e_i, e_i) = \mathbb{E}(\phi''(f'\xi, \xi)).$$


$$F(\text{diagram 4}) = \sum_i \phi''(e_i, e_i) = \Delta\phi = \mathbb{E}(\phi''(\xi, \xi)).$$


$$F(\text{diagram 5}) = \sum_{i,j} \phi''(e_i, f'''(e_j, e_j, e_i)) = \sum_i \phi''(e_i, (\Delta f)'(e_i)).$$


Related algebraic structures: E. Bronasco, Exotic B-series and S-series, 2024.

Study of **affine equivariant property** by McLachlan, Modin, Munthe-Kaas, Verdier, 2016  
and **orthogonal equivariant maps** by Laurent, Munthe-Kaas, 2023.

# Integration by parts using exotic aromatic trees: example

$$\begin{aligned}
 \int_{\mathbb{R}^d} F(\text{diagram}) (\phi) \rho_\infty dy &= \int_{\mathbb{R}^d} (\Delta \phi)' f \phi \rho_\infty dy = \sum_{i,j} \int_{\mathbb{R}^d} \frac{\partial^3 \phi}{\partial x_i \partial x_j \partial x_j} f_i \rho_\infty dy \\
 &= - \sum_{i,j} \left[ \int_{\mathbb{R}^d} \frac{\partial \phi}{\partial x_i \partial x_j} \frac{\partial f_i}{\partial x_j} \rho_\infty dy + \int_{\mathbb{R}^d} \frac{\partial \phi}{\partial x_i \partial x_j} f_i \frac{\partial \rho_\infty}{\partial x_j} dy \right] \\
 &= - \int_{\mathbb{R}^d} F(\text{diagram}) (\phi) \rho_\infty dy - \int_{\mathbb{R}^d} F(\text{diagram}) (\phi) \rho_\infty dy.
 \end{aligned}$$

where we used  $\nabla \rho_\infty = f \rho_\infty$ . We obtain: .

# Convergence analysis in the variable diffusion case

New notation for **exotic aromatic trees**:  $EAT = \{\bullet, \circlearrowleft, \circlearrowright, \dots\} = \{\bullet, \textcircled{1}, \textcircled{1}, \bullet, \textcircled{1}, \textcircled{1}, \dots\}$

$$\mathcal{L}\phi = \phi'F + \frac{\sigma^2}{2} \sum_{i=1}^d \phi''(D_i, D_i) = \mathcal{F}\left(\bullet + \frac{1}{2}\circlearrowleft\right)[\phi] = \mathcal{F}\left(\bullet + \frac{1}{2}\textcircled{1}\textcircled{1}\right)[\phi]$$

$$\mathcal{L}^2 = \mathcal{F}\left(\bullet\bullet + \bullet + \bullet\textcircled{1}\textcircled{1} + \textcircled{1}\textcircled{1} + \textcircled{1}\bullet + \frac{1}{2}\textcircled{1}\textcircled{1} + \frac{1}{4}\textcircled{1}\textcircled{1}\textcircled{2}\textcircled{2} + \textcircled{1}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{2}\right)$$

**Theorem (Bronasco, 2024, Bronasco, Leimkuhler, Phillips, and V., submitted 2025)**

We can use **integration by parts** denoted by  $\sim$  to modify  $\mathcal{A}_k$  without changing the value of  $A_k^* \rho_\infty$ . The order  $p$  condition becomes

$$(a \circ A)(\tau) = 0, \quad \text{for all } \tau \in EAT, |\tau| \leq p,$$

where  $A$  is an adjoint operation of the integration by parts.

# Convergence analysis in the variable diffusion case

New notation for **exotic aromatic trees**:  $EAT = \{\bullet, \circlearrowleft, \circlearrowright, \dots\} = \{\bullet, \textcircled{1}\textcircled{1}, \bullet\textcircled{1}\textcircled{1}, \dots\}$

$$\mathcal{L}\phi = \phi'F + \frac{\sigma^2}{2} \sum_{i=1}^d \phi''(D_i, D_i) = \mathcal{F}\left(\bullet + \frac{1}{2}\circlearrowleft\right)[\phi] = \mathcal{F}\left(\bullet + \frac{1}{2}\textcircled{1}\textcircled{1}\right)[\phi]$$

$$\mathcal{L}^2 = \mathcal{F}\left(\bullet\bullet + \textcolor{red}{\bullet} + \bullet\textcircled{1}\textcircled{1} + \textcircled{1}\textcircled{1} + \textcircled{1}\bullet + \frac{1}{2}\textcircled{1}\textcircled{1} + \frac{1}{4}\textcircled{1}\textcircled{1}\textcircled{2}\textcircled{2} + \textcircled{1}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{1}\textcircled{1}\right)$$

$$2\mathcal{A}_2 = \mathcal{F}\left(\bullet\bullet + \bullet\textcircled{1}\textcircled{1} + \frac{1}{2}\textcolor{red}{\textcircled{1}}\textcolor{red}{\textcircled{1}} + \textcircled{1}\bullet + \frac{1}{8}\textcolor{red}{\textcircled{1}}\textcolor{red}{\textcircled{1}} + \frac{1}{4}\textcircled{1}\textcircled{1}\textcircled{2}\textcircled{2} + \textcircled{1}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{2} + \frac{1}{2}\textcircled{2}\textcircled{1}\textcircled{1}\right)$$

## List of order conditions for the variable diffusion case

There are 93 order conditions for order 2 for the invariant measure sampling!

$$1. a(\bullet \textcircled{1} \textcircled{1}) - 2a(\textcircled{1} \textcircled{1} \textcircled{2} \textcircled{2}) = 0,$$

$$2. a(\textcircled{1} \textcircled{2} \textcircled{2}) - 2a(\textcircled{1} \textcircled{1} \textcircled{2} \textcircled{2}) = 0,$$

$$3. a(\textcircled{1} \textcircled{2} \textcircled{2}) - 2a(\textcircled{1} \textcircled{1} \textcircled{2} \textcircled{2}) = 0,$$

$$4. a(\textcircled{1} \textcircled{1}) = 0,$$

⋮

$$93. a(\textcircled{1} \textcircled{\hspace{0.5cm}}) = 0.$$

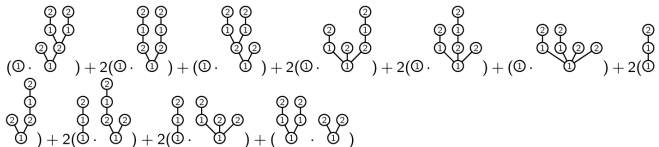
```
> f1 = [PT 1 [PT 2 []], PT 1 [PT 2 []]]
> display $ vector f1
```



```
> f2 = [PT 1 [], PT 1 [PT 2 [], PT 2 []]]
> display $ vector f2
```



```
> display $ graftFF f1 f2
```



**Ongoing work:** design of a systematic symbolic manipulation package in Haskell for exotic trees and forests (with E. Bronasco and J.L. Falcone).

# Conclusion

- preconditioning for enhancing the convergence to equilibrium of ergodic dynamics.
- postprocessing allows to increase the order of accuracy for sampling the invariant measure without increasing the number of force evaluations (same as the simplest Euler-Maruyama method).
- construction in the stochastic context inspired by [geometric numerical integration](#).

## Related ongoing work:

- preconditioned integrators for SPDEs, stochastic heat equation (with A. Debussche, C.-E. Bréhier, and A. Laurent).
- design of a systematic symbolic manipulation package in Haskell for exotic trees and forests (with E. Bronasco and J.L. Falcone).
- accelerating the convergence to equilibrium using reversible perturbations (Stratonovitch noise), with G. Pavliotis.

[Papers and preprints](http://www.unige.ch/~vilmart) available at: [www.unige.ch/~vilmart](http://www.unige.ch/~vilmart)

## Application: high order integrator based on modified equations

It is possible to construct integrators of weak order 1 that have order  $p$  for the invariant measure.

### Theorem (Abdulle, V., Zygalakis, 2015)

Consider an ergodic integrator  $X_n \mapsto X_{n+1}$  (with weak order  $\geq 1$ ) for an ergodic SDE in the torus  $\mathbb{T}^d$  (with technical assumptions),

$$dX = f(X)dt + g(X)dW.$$

Then, for all  $p \geq 1$ , there exists a modified equation

$$dX = (f + hf_1 + \dots + h^{p-1}f_{p-1})(X)dt + g(X)dW,$$

such that the integrator applied to this modified equation has order  $p$  for the invariant measure of the original system  $dX = fdt + g dW$  (assuming ergodicity).

Related work on modified equations for SDEs: Shardlow (2006, strong), Zygalakis, (2011, weak), Debussche & Faou, (2011, ergodic problems), Abdulle Cohen, V., Zygalakis (2013, weak), Bronasco, Laurent (2024, ergodic, algebraic structures).



## Example of high order integrator for the invariant measure (Brownian dynamics)

### Theorem (Abdulle, V., Zygalakis, 2015)

For  $p \geq 1$  and Brownian dynamics  $dX = f(X)dt + \sigma dW$ ,  $f = -\nabla V$ , the Euler-Maruyama scheme  $X_{n+1} = X_n + hf(X_n) + \sigma \Delta W_n$  applied to the modified SDE

$$dX = (f + hf_1 + h^2 f_2 + \dots + h^{p-1} f_{p-1})(X)dt + \sigma dW$$

$$f_1 = -\frac{1}{2}f'f - \frac{\sigma^2}{4}\Delta f,$$

$$f_2 = -\frac{1}{2}f'f'f - \frac{1}{6}f''(f, f) - \frac{1}{3}\sigma^2 \sum_{i=1}^d f''(e_i, f'e_i) - \frac{\sigma^2}{4}f'\Delta f - \frac{\sigma^4}{6}(\Delta f)'f - \frac{\sigma^4}{24}\Delta^2 f,$$

...

has order  $p$  for the invariant measure of  $dX = f(X)dt + \sigma dW$  (assuming ergodicity).

Remark 1: the weak order of accuracy is only 1.

Remark 2: derivative free versions can also be constructed.

Related: algebraic structures based on exotic aromatic trees and forests (Laurent, V., 2020, Bronasco, 2024, Bronasco, Laurent, 2024).